# Astro 260 <br> Group Problems \#9 <br> 2005-April-16 

The Flat Robertson-Walker Metric:

$$
\begin{gathered}
d s^{2}=-d t^{2}+a^{2}(t)\left(d x^{2}+d y^{2}+d z^{2}\right) \\
d s^{2}=-d t^{2}+a^{2}(t)\left(d r^{2}+r^{2} d \theta^{2}+r^{2} \sin ^{2} \theta d \phi^{2}\right)
\end{gathered}
$$

1. In this problem, you will develop expressions for the coordinate distance and the lookback time to an object who is at comoving distance $r$ away from us.
(a) Consider a flat Universe which has two components in it: cold matter and vacuum energy. The density (in units of critical density) of these two components today are $\Omega_{\mathrm{M}}$ and $\Omega_{\Lambda}$. (The book uses $\Omega_{v}$ for vacuum, but $\Omega_{\Lambda}$ is slightly more standard notation.) The critical density is

$$
\rho_{c}=\frac{3 H_{0}^{2}}{8 \pi}
$$

where $H_{0}$ is the present expansion rate of the Universe.
Write down $\rho_{\mathrm{M}}$ and $\rho_{\Lambda}$ (which are both potentially functions of time) in terms of $\Omega_{\mathrm{M}}, \Omega_{\Lambda}, H_{0}$, and $a$ (the scale factor).
(b) Consider the Friedmann Eqation for the flat Universe:

$$
\dot{a}^{2}-\frac{8 \pi \rho}{3} a^{2}=0
$$

Rewrite the Friedmann equation as a differential equation between $z$ and $t$, where $z$ is the cosmological redshift: $(1+z)=a_{0} / a(z)$, where $a_{0}=a\left(t_{0}\right)$ is the scale factor today (assume $a_{0}=1$ ), and $a(z)$ is the scale factor when light was emitted by an object that we observe to have redshift $z$.
(c) Show that the lookback time to an object at redshift $z$ (that is, the amount of Robertson-Walker time $t$ that has elapsed in between the emission and detection of light from an object whose redshift we observe to be $z$ ) has the value:

$$
t_{0}-t_{\mathrm{emit}}=\int_{0}^{z} \frac{-d z}{H_{0}(1+z)\left[\Omega_{\Lambda}+\Omega_{\mathrm{M}}(1+z)^{3}\right]^{1 / 2}}
$$

I hope you can understand my sloppiness with $z$; if not, then replace $z$ with $z^{\prime}$ everywhere inside the integral, and replace $d z$ with $d z^{\prime}$, and then it all makes sense.
(d) Write down a similar expression for $r$, the coordinate distance to that object. (Hint: consider $d r / d t$ for a photon travelling from there to here.)
(e) What is the physical distance from us to the emitting galaxy holding $t$ constant at $t=t_{0}$ ?
(f) What is the physical distance from us to the emitting galaxy holding $t$ constant at $t=t_{\text {emit }}$ ? (Your answer should be in terms of $z, H_{0}, \Omega_{\mathrm{M}}$, and $\Omega_{\Lambda}$.)
2. Cosmological Time Dilation: Consider the situation below:


A distant galaxy of measured redshift $z$ emits a photon at time $t_{\text {emit }}+d t_{\text {emit }}$ (when it has redshift $z+d z_{\text {emit }}$ ), and another one at time $t_{\text {emit }}$ (when it has redshift $z$ ). We detect the first photon at time $t_{0}+d t_{0}$ (note: both $d t_{\text {emit }}$ and $d t_{0}$ are negative), and the second photon now (at time $t_{0}$ ). We can be said to have had a redshift $d z_{0}$ at the detection of the first photon, as the scale factor was slightly below 1 then, as shown.
(a) Write down an expression for $d z_{\mathrm{emit}} / d z_{0}$ in terms fo $H_{0}, z, \Omega_{\mathrm{M}}$, and $\Omega_{\Lambda}$. This should not be an integral expression! Hint: both photons travel the same coordinate distance! You developed that expression for the second photon in the previous problem. Hint 2: Remember that:

$$
\int_{x_{0}}^{x_{0}+\Delta x} f(x) d x \simeq f\left(x_{0}\right) \Delta x
$$

for small $\Delta x$ small.
(b) Develop an expression for $d t_{0} / d t_{\text {emit }}$ in terms of $H_{0}, z, \Omega_{\mathrm{M}}$, and $\Omega_{\Lambda}$. This should not be an integral expression! This problem may be a little tricky. Hint: If you can work out an expression in terms of $d z_{0}$ and $d z_{\text {emit }}$, you can use part (a) to get rid of one of the two. You can also work out an expression just between $d z_{0}$ and $d t_{0}$, if you think about it.
3. Luminosity Distance: We define luminosity distance $d_{L}$ such that:

$$
F=\frac{L}{4 \pi d_{L}{ }^{2}}
$$

where $L$ is the luminosity of an emitting object, and $F$ is the detected flux (energy per time per area). In a euclidean geometry where nothing is expanding, $d_{L}$ is exactly the same as regular distance.
(a) If $L$ is the luminosity of the emitting object, what is $d N / d t^{\prime}$, the number of photons emitted by that object in its own proper time $d t^{\prime}$ ? (For simplicitly, assume all the photons are of the same frequency.)
(b) If the object is at coordinate distance $r$, what is $A$, the area of the sphere over which the photons emitted at the same time as the ones we detect from the object are spread?
(c) Write down an expression for $\frac{d N}{d A d t}$, where $d N$ is the number of photons emitted by a distance source in its proper time $d t^{\prime}$ and $d t$ is the amount of time it takes us to detect those photons. Your expression should be in terms of $L, \omega^{\prime}$ (the frequency of the emitted photons), and $r$.
(d) Write down an expression for $\frac{d E}{d A d t}$, the flux we detect, in terms of the quantities from (c) and $\omega$, the detected photon frequency.
(e) Convert your expression in (d) to an expression in terms of $z, \Omega_{\mathrm{M}}, \Omega_{\Lambda}$, and $H_{0}$. This is the standard luminosity distance formula. By measuring $F$ and $z$ for a lot of objects of known $L$, we can fit the values of the parameters $H_{0}, \Omega_{\mathrm{M}}$, and $\Omega_{\Lambda}$ to this curve.

